

Manual for PeTeR

Version 2.1.0

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1 Introduction

The flexible C++ program PeTeR [1] was developed for the purpose of performing the expensive numerical computation of the resummed cross section according to [2, 3, 4], i.e. the transverse momentum distribution of electroweak boson production at hadron colliders (γ , Z , W^\pm , H). It consists of two parts, the Unix-like command-line program `peter`, which provides a user-friendly interface, and `libpeter`, which contains all the computational ingredients. `libpeter` is a library, and as such it may easily be used to build other programs.

So far, PeTeR is able to calculate the NLL, N²LL, N³LL resummed cross sections as well as the corresponding singular ones, LO, NLO_{sing} and NNLO_{sing}. In addition, the full NLO fixed-order results for vector boson and for Higgs production are implemented and were extensively cross-checked against QT [5, 6]¹ and HqT2.0 [7].

PeTeR includes the full two-loop hard functions for electroweak boson production computed in [4], but for simplicity the tiny two-loop contributions which arise when a neutral vector boson couples to an internal quark loop² are neglected (N_V^v and N_V^a , cf. Eqs.(5.2) and (5.3) of [4]).

2 Installation

The code can be downloaded from <http://peter.hepforge.org/>. The gzipped tar file unpacks into a directory PeTeR-k.m.n, where k.m.n denotes the version number. Change into this directory and type

```
./configure
make
```

A successful compilation creates `src/libpeter/libpeter.a` and `src/peter/peter`. It is possible to run the code directly inside `src/peter`. In order to install/copy `peter` and `libpeter` into `/usr/local` type

¹I am grateful for discussions with Richard Gonsalves.

²The axial part at two-loop level is still unknown. It is, however, already very small at one loop.

```
make install
```

Super user privileges may be necessary (type `sudo make install`). The directory can be changed by the option `--prefix=directory` of the `configure` command.

If the library LHAPDF [8], version 5.4.0 or higher (but still version 5, version 6 is not yet tested), is installed, its capabilities can be included by specifying the configure option `--with-LHAPDF`. If the LHAPDF library is installed in a non-standard directory, i.e. the executable `lhpdf-config` is not in a directory included in the environment variable `$PATH`, one has to specify the complete path to the executable `lhpdf-config`

```
./configure --with-LHAPDF \  
LHAPDF_CONFIG=/path-to-lhpdf-config/lhpdf-config
```

How to choose a pdf from LHAPDF in the executable `peter` is explained in the next chapter. More details about the installation process are given in the file `INSTALL` and by invoking `./configure --help`.

3 Command Line `peter`

In order to facilitate the user interface of the command line program `peter`, the whole command line parsing is done by the `boost/program_options` library [9]. After storing all input parameters, `peter` calls the function `sigma_pT` of `libpeter`. Note that `peter` automatically takes advantage of parallelization on multi-core processors by the build-in functionality of the `CUBA` library [10]. An overview of library dependencies is depicted in Figure 1. A simple call `peter -h` produces the following detailed help message:

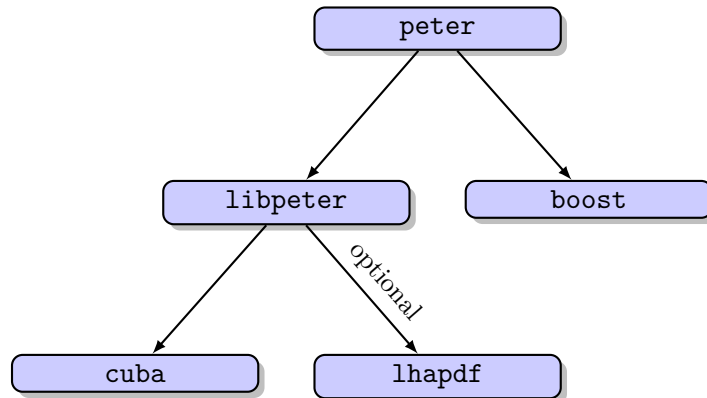


Figure 1: Dependencies of `peter`

Call: `peter [options]`

Description: Specify at least one option otherwise `'-h'` is assumed. `Peter` calculates the transverse momentum distribution of electroweak boson production at colliders such as the LHC: `dsigma/dpT[pbarn]` of `hadron + hadron -> V + X`.

See <http://peter.hepforge.org/> for more informations.

Allowed options:

-h [--help]	Print help message.
--version	Print version number.
-v [--verbosity] arg (=1)	Set the level of verbosity: 0=silent 1=input values (reusable as input file by copy paste) 2=add pdf information if available (not yet working for LHAPDF)
-f [--configfile] arg	Set the name for a configuration file that will be parsed. Command line input overwrites config file input.
--boson arg (=1)	Set the boson being produced: 0 = Photon 1 = Z 2 = W+ 3 = W- 4 = Higgs
--collider arg (=0)	Set the collider: 0 = proton-proton 1 = proton-antiproton 2 = antiproton-antiproton
--cms-energy arg (=8000)	Set cms-energy sqrt(S) [GeV].
-p [--transverse-momentum] arg (=100)	Set pT start point [GeV].
--pT-steps arg (=1)	Set the number of pT points to be calculated.
--pT-step-size arg (=10)	Set the step size between pT points [GeV].
--log-scale [=arg(=1)] (=0)	If true, set logarithmic steps. i=0..n-1, n=pTsteps, d=pT-step-size linear: pTstep(i) = pT + i*d, i=0..n-1 log: pTstep(i) = pT * ((pT+(n-1)*d)/pT)^(i/(n-1)) So both cases start with pT and end with pT + (n-1)*d.
--min-rapidity arg (=-1000000)	Set minimal rapidity.
--max-rapidity arg (=1000000)	Set maximal rapidity.
--hard-scale arg (=-1)	Set hard scale muH[GeV]. The general scale choices are >0 : mu = input [GeV] -1 : mu = (13 pT + 2 M)/12 - pT^2/sqrt(S) -2 : mu = (7 pT + 2 M)/12 * (1 - 2 pT/sqrt(S)) -3 : mu = mu_jet^2/mu_hard -4 : mu = sqrt(pT^2 + M^2)

where M is the boson mass.
Option -3 is forbidden for the hard scale.

--jet-scale arg (=-2) Set jet scale μ_J [GeV]. Same options as hard scale. Option -3 is forbidden.

--soft-scale arg (=-3) Set soft scale μ_S [GeV]. Same options as hard scale.

--factorization-scale arg (=-1) Set factorization scale μ_F [GeV]. Same options as hard scale. If jet, soft and factorization scale are equal, η equals 0. For each partonic channel, the singularity $1/\Gamma[\eta]$ is treated by expanding in η for $|\eta| < 0.001$.

--hard-scale-factor arg (=1) Multiply the hard scale.

--jet-scale-factor arg (=1) Multiply the jet scale.

--soft-scale-factor arg (=1) Multiply the soft scale.

--factorization-scale-factor arg (=1) Multiply the factorization scale.

--flavors arg (=5) Set the number of active flavors.

--order-resum arg (=0) Set the order of resummation:
3 = N^3LL
2 = NNLL
1 = NLL
0 = fixed-order LO (no resummation)
-1 = fixed-order NLO (no resummation)
-2 = matched NNLL+NLO
-3 = matched $N^3LL+NLO$

For resummed results: $h*j*s$, $\log(U)$, η are expanded up to $\alpha_s^{(order-1)}$. If all scales are set equal to μ_F one gets the singular/threshold expanded cross section N^xLO_{sing} with $x=order-1$.
For fixed-order results: $\mu_R=hard$ scale, $\mu_F=fact.$ scale.
For matched results:
 $resummed(\mu_H, \mu_J, \mu_S, \mu_F) + nlo(\mu_F) - nlo_sing(\mu_F)$.

Please note that the matched results (-1,-2,-3) for photon production are incomplete (!) since they do not include the NLO photon fragmentation contribution. To obtain finite results without fragmentation, the photon mass parameter needs to be kept at a non-zero value and acts as a collinear cutoff.

--order-alphas arg (=0) Set the order of the running of alphas:
0 = 1-loop (beta0)
1 = 2-loop (beta1)

$2 = 3\text{-loop} (\beta_2)$
 $3 = 4\text{-loop} (\beta_3)$.

`--alphas-epsrel arg (=1e-06)` Set the relative precision for the computation of the running of alphas. Newton's method quits with a warning after 1000 iterations if accuracy is not reached.

`--with-triangles [=arg(=1)] (=0)` Include triangle contributions for b and t quarks (only Z bosons). Quark masses $\geq 10^6$ GeV are treated as infinitely heavy.

`--with-two-loop-const [=arg(=1)] (=0)` Include the two loop constant in the hard function. Without it, the α_s^2 part of the hard function is normalized as $h_2(\mu_H=p_T)=0$.

`--with-ct2-evolution [=arg(=1)] (=1)` For Higgs only: If true, the top-quark matching scale $\mu_t=mass_t$. If false, $\mu_t=\mu_H$. For resummed and fixed-order results, the Wilson coefficient Ct_2 is evaluated at scale μ_t , the cross section is multiplied by the evolution factor $U_{Ct_2}(\mu_t, \mu_H)$. Note that $U_{Ct_2}(\mu_H, \mu_H)=1$. For strict fixed-order results, i.e. `order-resum=-1` (NLO), one has to set this option to false.

`--no-hard [=arg(=1)] (=0)` Switch off the hard function, i.e. sets hard functions to 1.

`--no-jet [=arg(=1)] (=0)` Switch off the jet function, i.e. sets jet functions to 1.

`--no-soft [=arg(=1)] (=0)` Switch off the soft function, i.e. sets soft functions to 1.

`--scale-variation arg (=0)` Compute theoretical errors from scale variation by a factor of 1/2 and 2:

- 0: Don't compute scale variation.
- 1: Add the errors in the additional columns `scale+` and `scale-`.
- 2: Additionally output the scale varied cross sections.

 For LO, NLO and singular: The scale $\mu_F=\mu_R$ is varied up and down. Then, a parabola is fitted through the 3 points $\mu_{var}/\mu=(1/2,1,2)$, and max and min values in this interval minus the central value ($\mu_{var}/\mu=1$) give the error estimates.

Resummed and matched: Same procedure, but each scale μ_H, μ_J, μ_S and μ_F is varied separately. The individual

errors are added in quadrature.
Note: The central value ($\mu_{\text{var}}/\mu=1$) might differ from the one computed without scale variation (within the given integration error), as the 3 and 9 cross sections, respectively, are computed at once. As the scale variation is usually larger than a few percent, one might consider setting `cuba-epsrel` not too small.

`--subtractions arg (=2)` Set the order of subtractions of the resummed integrand. 0-2 are implemented.

`--MH arg (=126)` Set Higgs boson mass [GeV].

`--MP arg (=1)` Set photon mass [GeV]. See option `order-resum`.

`--MZ arg (=91.1874)` Set Z boson mass [GeV].

`--MW arg (=80.381)` Set W boson mass [GeV].

`--sin2theta arg (=0.22296)` Set $\sin^2(\theta_w)$.

`--alpha arg (=0.00781592)` Set electroweak coupling constant α .

`--GF arg (=1.16638e-05)` Set Fermi constant G_F . It is used for Higgs production only (instead of α, \dots).

`--alphas arg (=0.11707)` Set strong coupling constant α_s at M_Z .

`--mu-alphas arg` Set the default scale for α_s . Default is M_Z .

`--Vud arg (=0.97427)` Set CKM matrix element V_{ud} .

`--Vus arg (=0.22534)` Set CKM matrix element V_{us} .

`--Vub arg (=0.00351)` Set CKM matrix element V_{ub} .

`--Vcd arg (=0.2252)` Set CKM matrix element V_{cd} .

`--Vcs arg (=0.97344)` Set CKM matrix element V_{cs} .

`--Vcb arg (=0.0412)` Set CKM matrix element V_{cb} .

`--Vtd arg (=0.00867)` Set CKM matrix element V_{td} .

`--Vts arg (=0.0404)` Set CKM matrix element V_{ts} .

`--Vtb arg (=0.999146)` Set CKM matrix element V_{tb} .

`--mass-b arg (=0)` Set the b-quark mass.

`--mass-t arg (=173.5)` Set the t-quark mass.

`--pdf-type arg (=0)` Set the pdf type:
-1 = fake pdfs ' $x*(1-x)$ '
0 = MSTW2008nnlo
1 = LHAPDF

`--pdf-filename arg` For LHAPDF, set the filename of the pdf including the full path.

`--pdf-member arg (=0)` For LHAPDF, set the member of the pdf (if it has different members).

`--cuba-routine arg (=3)` Set the integration routine of cuba:
0 = Vegas

	1 = Suave
	2 = Divonne
	3 = Cuhre
	For more informations see the documentation of the CUBA library.
--cuba-epsrel arg (=0.001)	Set the maximal relative error of integration.
--cuba-epsabs arg (=1e-20)	Set the maximal absolute error of integration.
--cuba-verbosity arg (=0)	Set the verbosity level of the cuba integration.
--cuba-seed arg (=0)	Set the seed for the pseudo-number generator and chooses the generator.
--cuba-level arg (=0)	Choses the random number generator: seed level generator 0 any Sobol q !=0 0 Mersenne-Twister p !=0 !=0 Ranlux p q=quasi-random, p=pseudo-random
--cuba-mineval arg (=0)	Set minimum number of integrand evaluations required.
--cuba-maxeval arg (=100000)	Set the (approximate) maximum number of integrand evaluations allowed.
--cuba-last-sampling arg (=0)	It true, only the last (largest) set of samples is used in the final result.
--cuba-smoothing arg (=1)	(Vegas and Suave only) If true, apply additional smoothing to the importance function, this moderately improves convergence for many integrands.
--cuba-nnew arg (=2000)	(Suave only) Set the number of new integrand evaluations in each subdivision (approximate).
--cuba-flatness arg (=25)	(Suave only) Set a flatness parameter. The flater the higher the value might be.
--cuba-key0 arg (=9)	(Cuhre only) Choses the basic integration rule. key0=7,9,11,13 selects the cubature rule of degree key0. Note that the degree-11 rule is available only in 3 dimensions, the degree-13 rule only in 2 dimensions. For other values, the default rule is taken: degree-13 in 2 dimensions, degree-11 in in 3, degree-9 otherwise. The pT-spectrum of on-shell bosons has a 3 dimensional integration (y,x,mx2).
--cuba-key1 arg (=47)	(Divonne only) Determine sampling in the partition phase.
--cuba-key2 arg (=1)	(Divonne only) Determine sampling in

	the final integration phase.
--cuba-key3 arg (=1)	(Divonne only) Set the strategy for the refined phase.
--cuba-maxpass arg (=5)	(Divonne only) Control the thoroughness of the partitioning phase integration phase.
--cuba-border arg (=1e-10)	(Divonne only) Set the width of the border of the integration region. Use a non-zero border if the integrand subroutine cannot produce values directly on the integration boundary.
--cuba-maxchisq arg (=10)	(Divonne only) Set the maximum χ^2 value a single subregion is allowed to have in the final integration phase.
--cuba-mindeviation arg (=0.25)	(Divonne only) Set a bound, given as the fraction of the requested error of the entire integral, which determines whether it is worthwhile further examining a region that failed the χ^2 test. Only if the two sampling averages obtained for the region differ by more than this bound is the region further treated.

Options can be specified on the command line and by providing a configuration file via `--configfile filename` (or `-f filename`). Note that the leading ‘--’ of the option names must be omitted in configuration files and that command line input overwrites configuration file input. If you call `peter` with `--verbosity 1` (or `-v1`), it prints the options in a format suitable for copy & paste into a configuration file. In addition, the file `sampleInput.txt` in `src/peter` is a good starting point for your own configuration file. The call

```
./peter -v0 -f sampleInput.txt --boson 4 --cms-energy 8000 \
--order-resum 2 -p 50 --pT-steps 4 --pT-step-size 50 --cuba-routine 2
```

produces an output that looks like

pT [GeV]	dsigma/dpT [pbarn/GeV]	error [pbarn/GeV]	prob	fail
50	0.0505339	5.06951e-05	4.10783e-06	0
100	0.0100421	1.00097e-05	1.67229e-08	0
150	0.0030726	3.1214e-06	1.75918e-05	0
200	0.0011569	1.13585e-06	7.11653e-14	0

The `error` is the presumed absolute error of integration, `prob` is the χ^2 -probability that `error` is not a reliable error estimate, and `fail` evaluates to 0 if the desired accuracy was reached; for more details see [10].

If you compiled PeTeR with LHAPDF, you can specify a pdf by the options `--pdf -type 1 --pdf-filename /path-to-share/lhapdf/PDFsets/xxx.LHgrid`, where `xxx`

is the pdf name of your choice. As you see, the option `pdf-filename` requires the full path of the pdf grid file.

As a last point, note that you can redirect the output of `peter` in the usual way, e.g. `peter 1> output.dat 2> error.log` redirects the standard output into the file `output.dat` and all warnings and error messages into `error.log`.

4 Library `libpeter`

Everything in the library `libpeter` is declared in the namespace `peter`. The main function in the file `sigma_pT.hpp`

```
std::vector<double> sigma_pT(const Cuba_Parameters& cp, const
    Process_Data& pd, const SM_Parameters& sm, const PDF& pdf)
```

provided by the library `libpeter` integrates the cross section via the CUBA library and returns the result as a vector in the form $(p_T, \frac{d\sigma}{dp_T}, \mathbf{error}, \mathbf{prob}, \mathbf{fail})$. The input is given by the self explaining classes as shown above. The structure of the library can be seen in [Figure 2](#). In the following, we give short summaries of the most important header files; the actual implementation might nevertheless be in the corresponding source file (`.cpp` instead of `.hpp`).

`alphas.hpp` The running of α_s up to 4-loop (β_3) is implemented by solving the equation

$$\ln \frac{\mu^2}{\mu_0^2} = \int_{\alpha_s(\mu_0)}^{\alpha_s(\mu)} d\alpha \left(-\alpha \sum_{n=0}^{\text{order}} \left(\frac{\alpha}{4\pi} \right)^{n+1} \beta_n \right)^{-1} \quad (1)$$

for $\alpha_s(\mu)$ with Newton's method. The integral is calculated analytically for the corresponding value of order (option `order-alphas`).

`mathfunctions.hpp` Several special functions are implemented at double precision, among these

- `pow10`: 10^x
- `dilog` and `cdilog`: dilogarithm $\text{Li}_2(x)$ for real and complex arguments
- `gamma`: Euler gamma function $\Gamma(x)$
- `polygamma`: polygamma functions $\Psi^{(n)}(x)$ for $n \in \{0, 1, 2, 3, 4, 5, 6\}$ (Higher n are implemented but with less accuracy since they are not required.)

All of the above functions have been extensively checked against their *Mathematica* implementations.

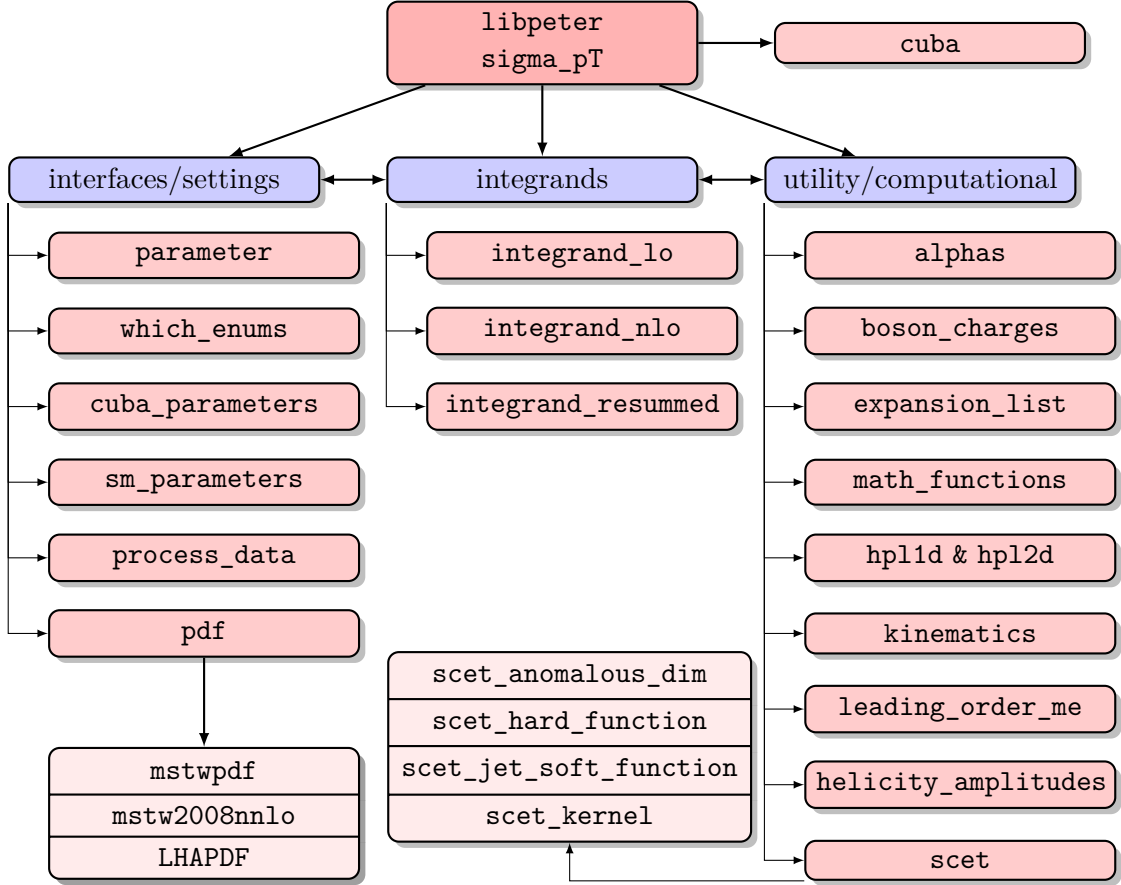


Figure 2: Schematic structure of `libpeter`. Blue boxes represent logical units, reddish boxes represent files (without indication for vector bosons `_V` and higgs boson `_H`, without file extensions `.hpp` and `.cpp`) or libraries.

hpl1d.hpp and hpl2d.hpp For the numerical computation of one- and two-dimensional harmonic polylogarithms, as required for the two-loop hard functions, these files provide full C++ versions of [11] and [12], which were originally written in `Fortran`.

The conversion has been achieved by modifying the `Fortran` code to comply to the C++ syntax. Apart from other differences between C++ and `Fortran`, the way of addressing elements of arrays was trivialized by template classes `Fortranarray1d`, `Fortranarray2d` and so forth. Schematically they look like

```
template <class T, const int n1, const int n2>
class FortranArray1d
{
private:
    const static int n_ = n2-n1+1;
    T value_[n_];
```

```

public:
    FortranArray1d() {
        for (int n=0; n<n_; ++n) value_[n] = T();
    };

    FortranArray1d(const T t[n_]) {
        for (int n=0; n<n_; ++n) value_[n] = t[n];
    };

    T operator() (int n) const {
        if (n<n1 || n2<n ) {
            /*error handling*/
        }
        return value_[n-n1];
    };

    T& operator() (int n) {
        if (n<n1 || n2<n ) {
            /*error handling*/
        }
        return value_[n-n1];
    };

    int size1() {return n_;};
};

```

Thus, no statement of accessing array elements on which the Fortran code heavily relies had to be changed. The evaluation of harmonic polylogarithms seems to be the bottleneck for $N^3\text{LL}$ and $\text{NNLO}_{\text{sing}}$ computations with the option `with-two-loop-const` enabled.

helicity_amplitudes.hpp The very long expressions for the helicity amplitudes necessary for NNLO hard functions are implemented by automatic insertion of the preprocessed expressions. It would be nice to implement the coefficients α , β , γ and δ of [13] as classes with a common interface. The easiest and most standard way of achieving this is to derive all of them from a common (base) class, which allows the storing of real and imaginary parts of coefficients of different orders in α_s , multiplying and complex conjugating them. This approach, however, has one drawback: it calls virtual functions, which are considered slow for high performance computations. Therefore, the *magic* of the Curiously Recursive Template Pattern (CRTP) [14] is invoked. There is a template class `template <class Derived> class Coefficient`, which provides the requested common interface. Every coefficient is then derived from a template of its own type, e.g. `class Alpha_2a : public Coefficient<Alpha_2a>`.

expansion_list.hpp In order to achieve the expansion of hard times jet times soft functions in α_s , these functions return the type `ExpList`, which is a class storing the

order. Actually, `ExpList` is just an alias for `ExpansionList<double>`, while `ExpansionList<T>` is a template class for any type `T`. The template function

```
template <class T>
ExpansionList<T> operator*(const ExpansionList<T>& left, const
    ExpansionList<T>& right)
{
    unsigned int order = left.order();
    if (right.order() < order) order = right.order();
    ExpansionList<T> result(order);
    result[0] = left[0] * right[0];
    if (order >= 1)
    {
        result[1] = left[0] * right[1] + left[1] * right[0];
    }
    if (order >= 2)
    {
        result[2] = left[0] * right[2] + left[2] * right[0] + left
            [1] * right[1];
    }
    return result;
};
```

provides the multiplication of two `ExpansionList<T>`s to the right order.

scet.hpp The header file `scet.hpp` simply includes `scet_anomalous_dim.hpp`, `scet_hard_function.hpp`, `scet_jet_soft_function.hpp` and `scet_kernel_function.hpp`. These implement hard, jet and soft functions, evolution factors, the kernel and its derivatives according to [2, 3, 4].

pdf.hpp There are two implementations for pdfs. One is an interface to the LHAPDF library [15, 8]. Another is a modified version of MSTW [16, 17], `mstwpdf.hpp`. The pdf `mstw2008nnlo` is hard coded in the C++ file `mstw2008nnlo.cpp`; neither an external data file nor an IO-operation is required anymore. The tradeoff in memory space is negligible for modern computers. An additional feature required for second subtractions is the possibility to compute the first derivative in x of pdfs.

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